

Neutron stars in 1+4D with interacting Fermi gas

Anna Horváth

*Wigner Research Centre for Physics
Eötvös Loránd University*

In collaboration with:

Gergely Gábor Barnaföldi

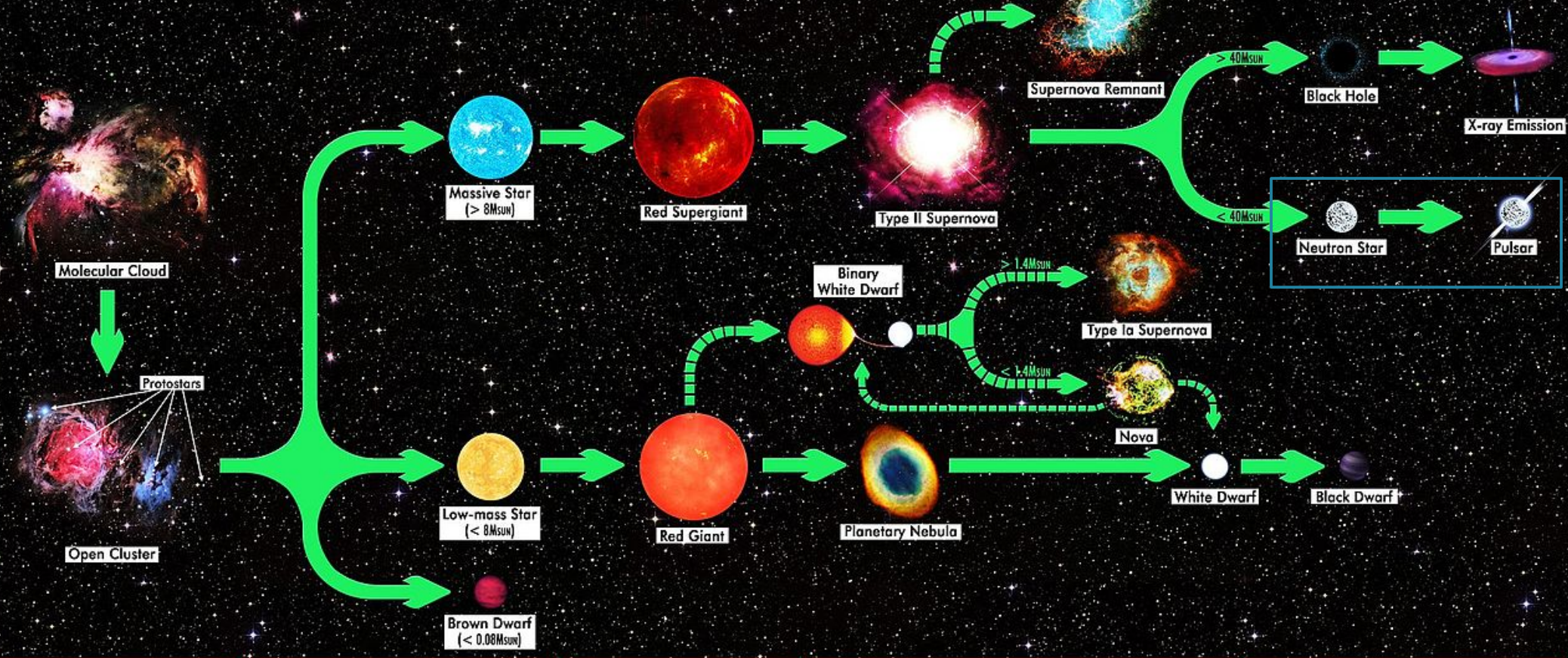
Wigner Research Centre for Physics

Emese Forgács-Dajka

Eötvös Loránd University



STELLAR LIFE CYCLE



Birth **Main Sequence** **Old Age** **Death** **Remnant**

Neutron stars

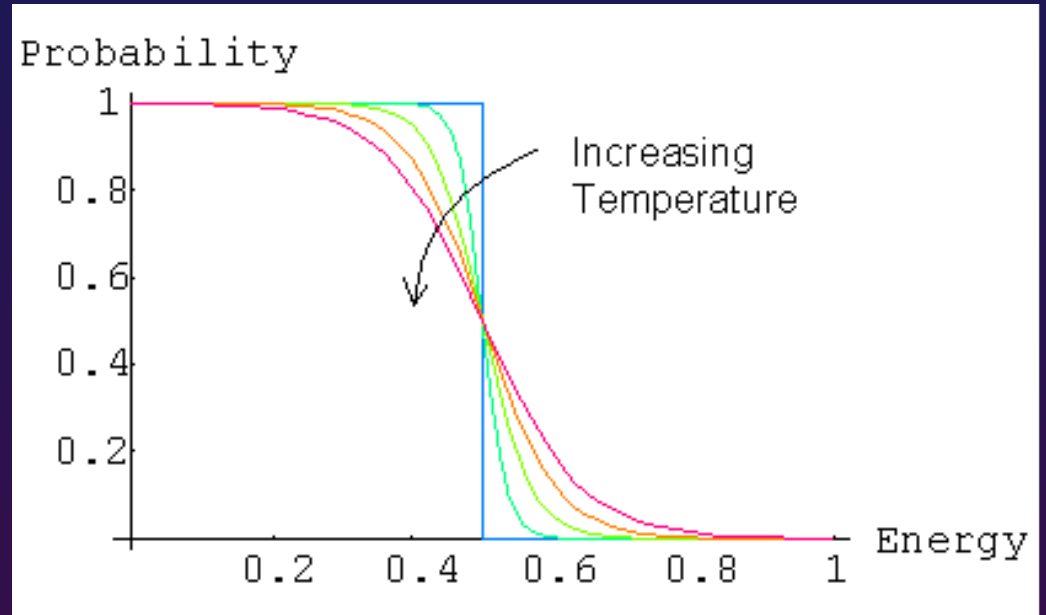
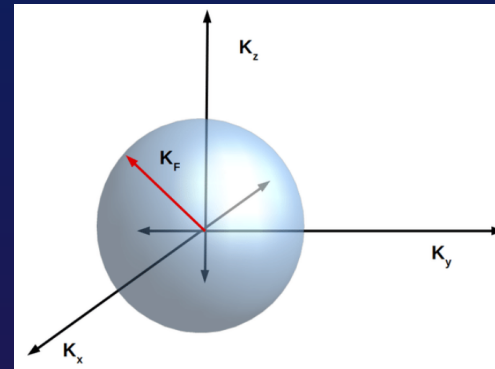
- Compact objects:

$$R \sim 10\text{km} \quad M \sim 1.1\text{--}2M_{\odot}$$

- Supported by **baryon degeneracy**
 - Pauli exclusion principle
- **Fermi-Dirac distribution**

$$\bar{n}_i = \frac{1}{e^{(\varepsilon_i - \mu)/k_B T} + 1}$$

- **Zero temperature approximation**



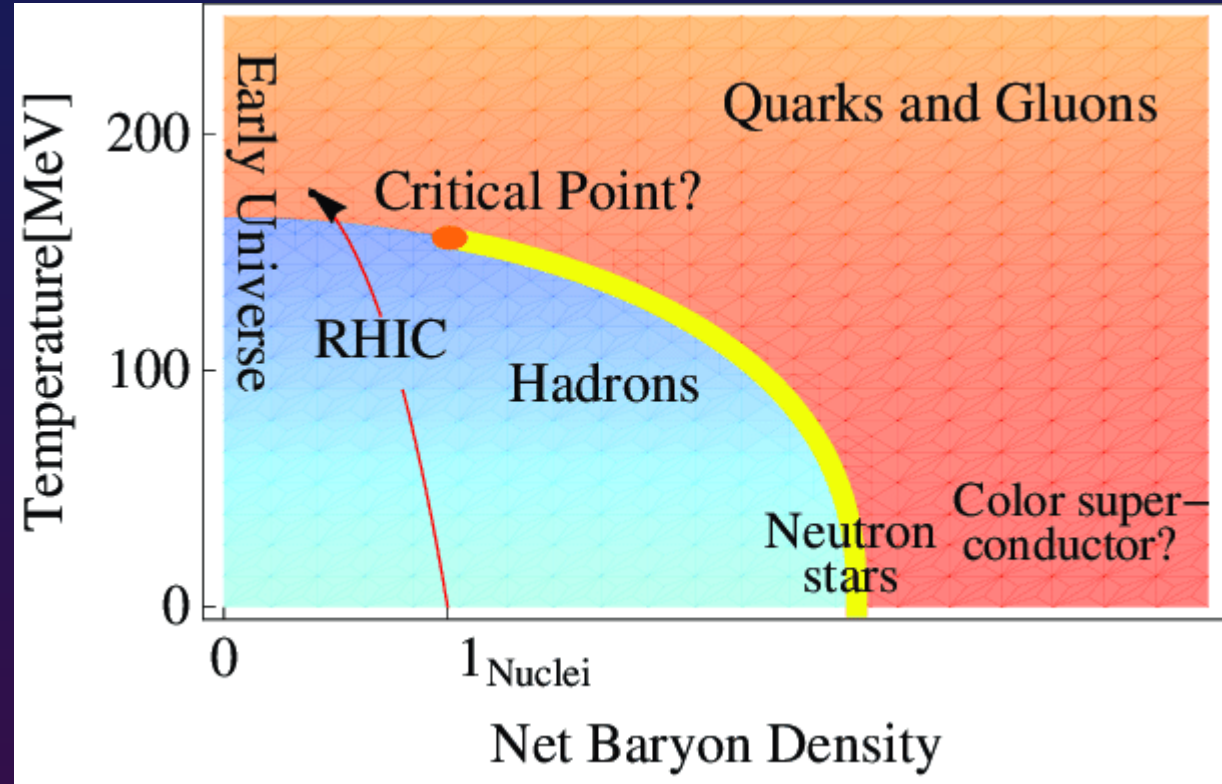
Fermi-Dirac distribution

QCD phase diagram

- Collider experiments probe high temperature, low density regime (RHIC, LHC)
- Neutron stars reside in the high density, low temperature regime

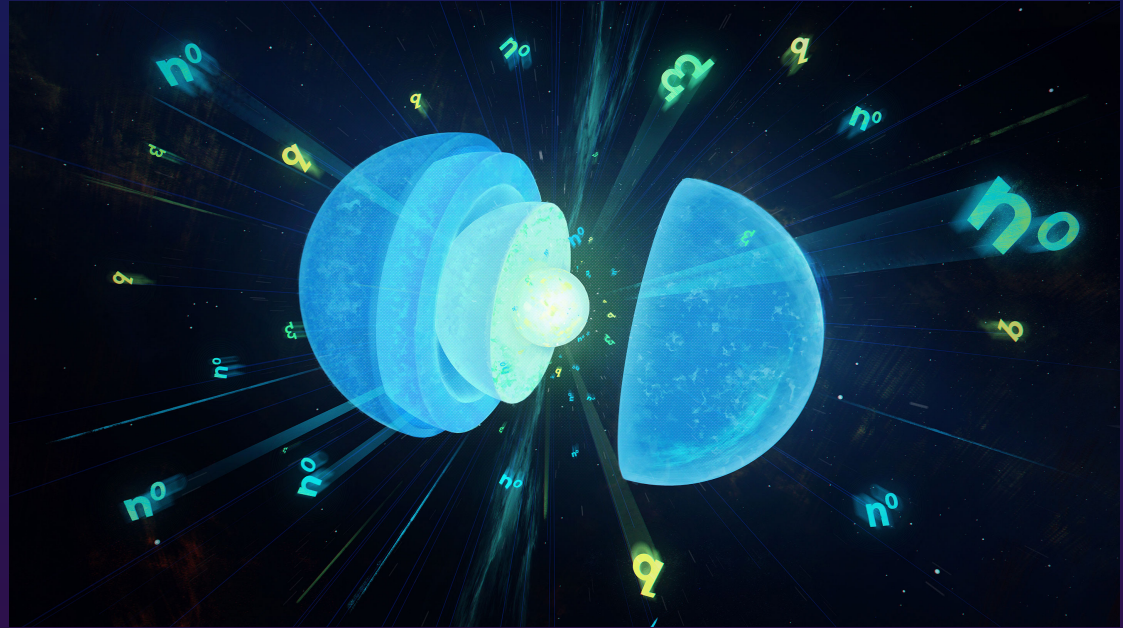


Probe physics in environments not available on Earth



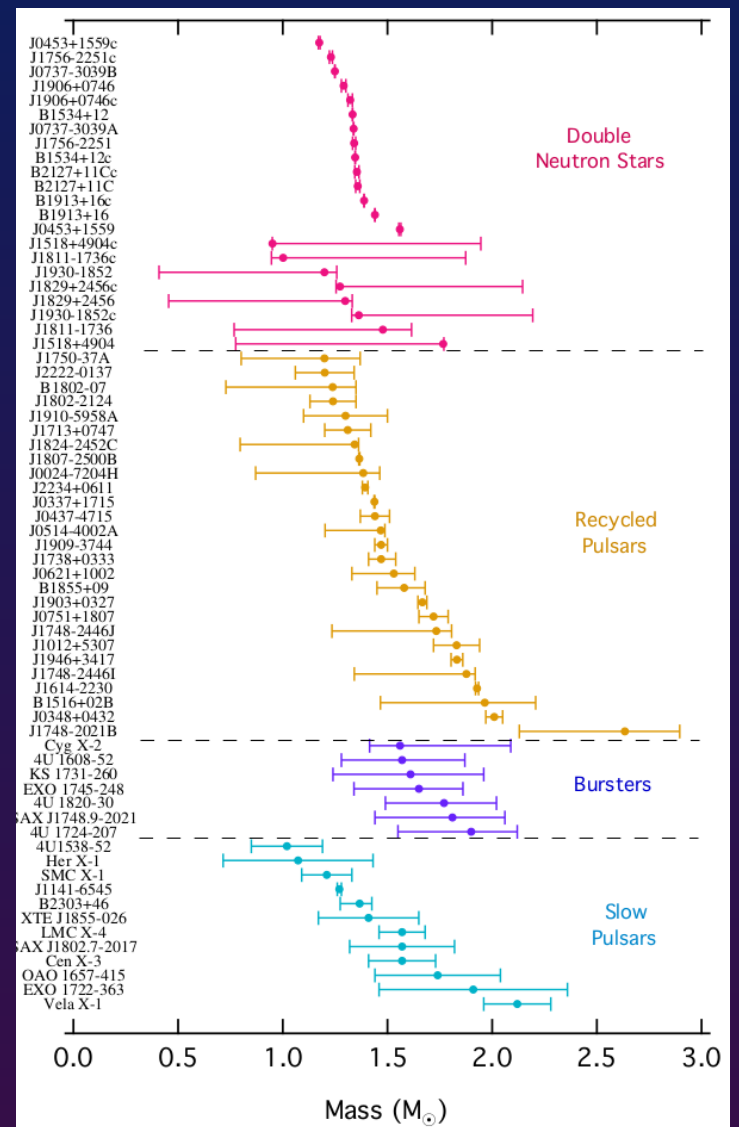
Probe physics of neutron stars

- Compare numerical simulations with experimental data – building up stars
- Main observables:
 - mass – difficult to measure
 - Error $\sim 2\text{--}5\%$
 - radius – even more difficult
 - Error $\sim 10\text{--}15\%$



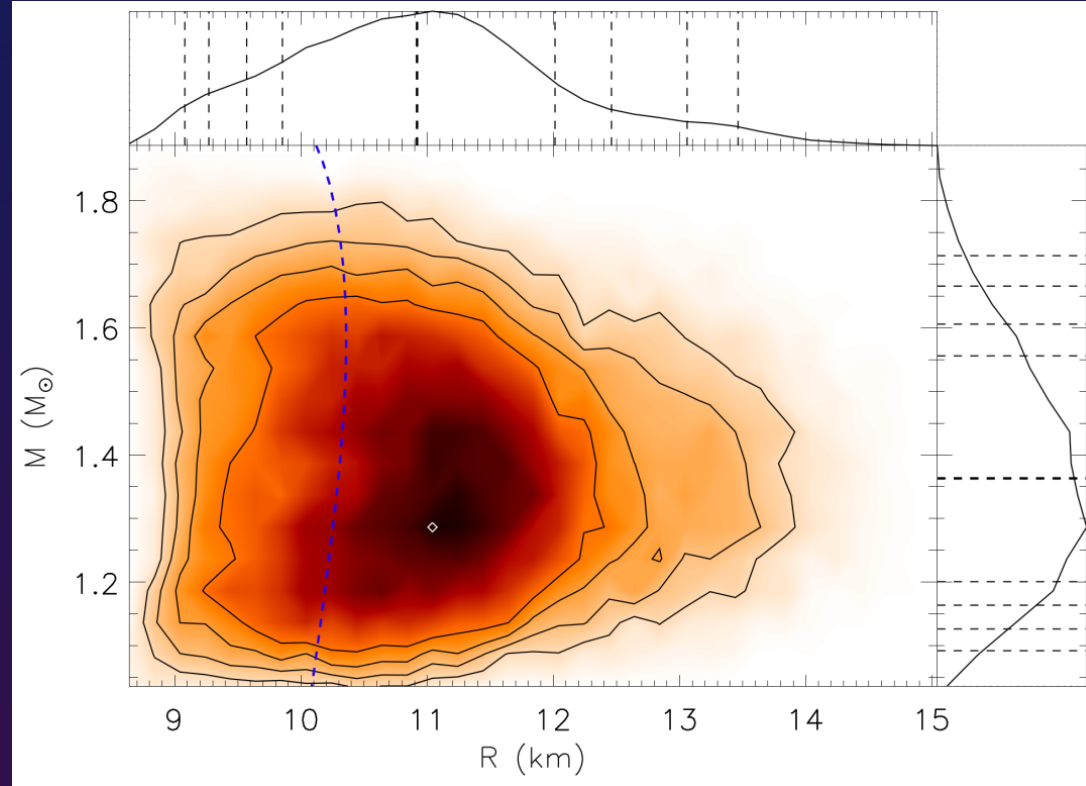
Mass measurements

- Radio measurement of rotation powered pulsars – most precise
- Also: X-rays, gamma rays, gravitational waves
- Binary or multiple component systems – tracking of orbital motion
- Better if both compact – point masses
- Isolated cannot be measured!
- Methods:
 - Advance of periastron
 - Einstein delay
 - Orbital period decay
 - Shapiro delay



Radius measurements

- Sources:
 - Nuclear physics
 - Astrophysical observations
 - GR assumptions
 - Causality of EoS
- Measurements based on **thermal emission** from surface:
 - Apparent angular size
 - Effects of NS spacetime on emission
 - **Spectroscopy** and **timing**
- **Gravitational waves** (figure: GW170817)



James M. Lattimer, “Neutron Star Mass and Radius Measurements” (2019)

Building up stars

- Two equations needed:
 - **Tolmann-Oppenheimer-Volkoff (TOV)** \sim GR hydrostatic equation

$$\frac{dp(r)}{dr} = -\frac{GM(r)\varepsilon(r)}{r^2} \times \left[1 + \frac{p(r)}{\varepsilon(r)}\right] \left[1 + \frac{4\pi r^3 p(r)}{M(r)}\right] \left[1 - \frac{GM(r)}{r}\right]^{-1}$$

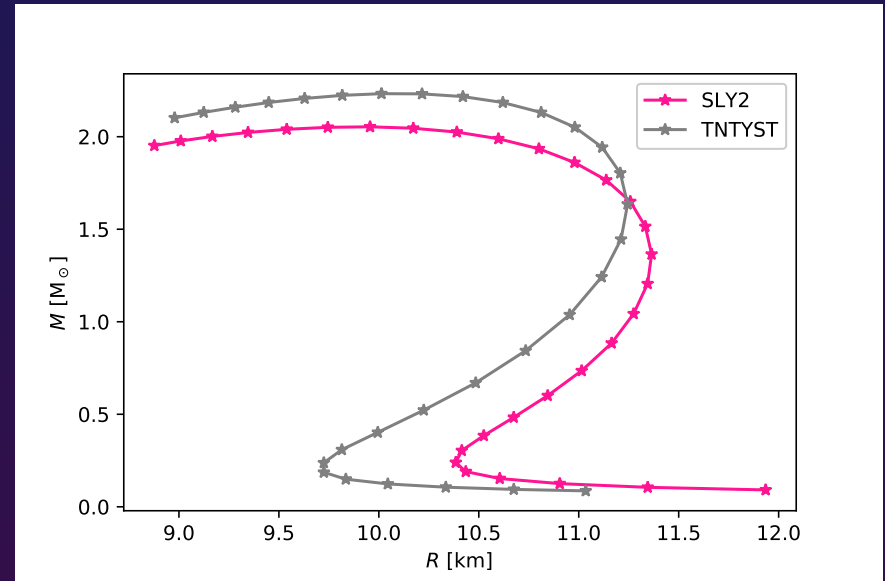
static,
spherical
symmetry

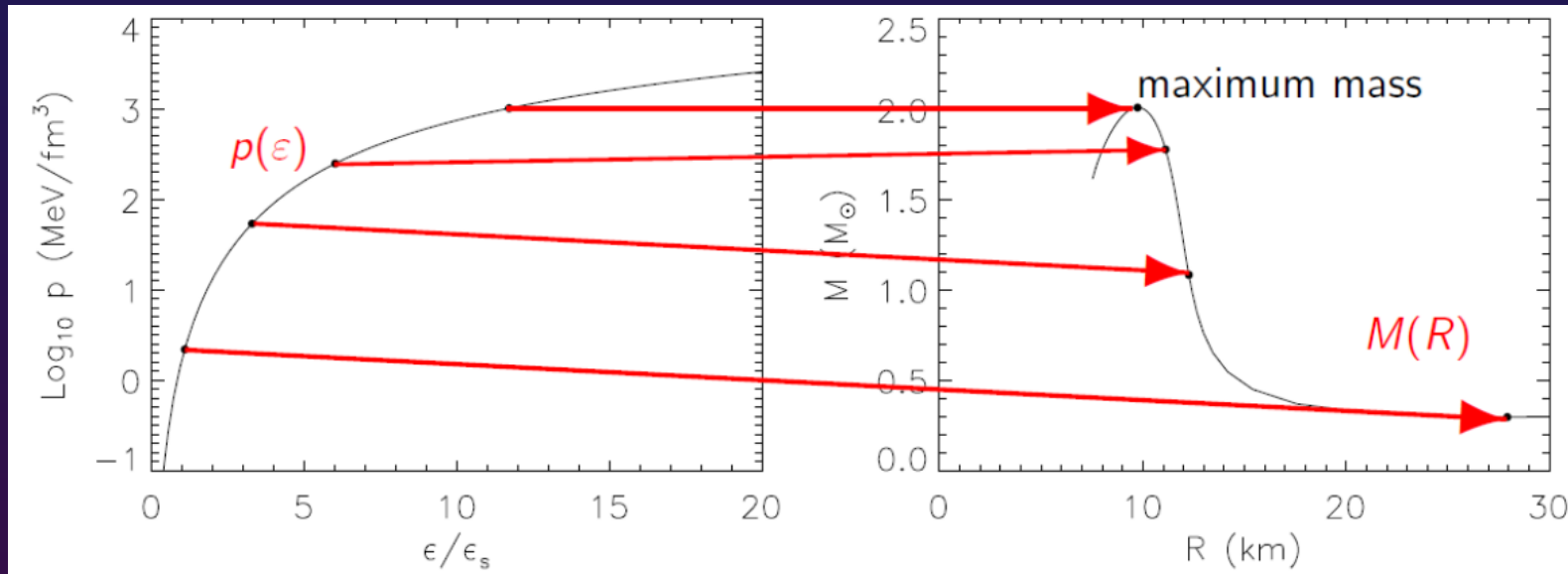
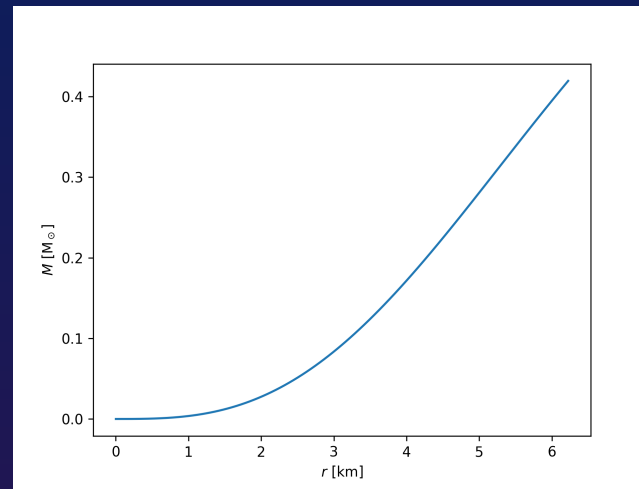
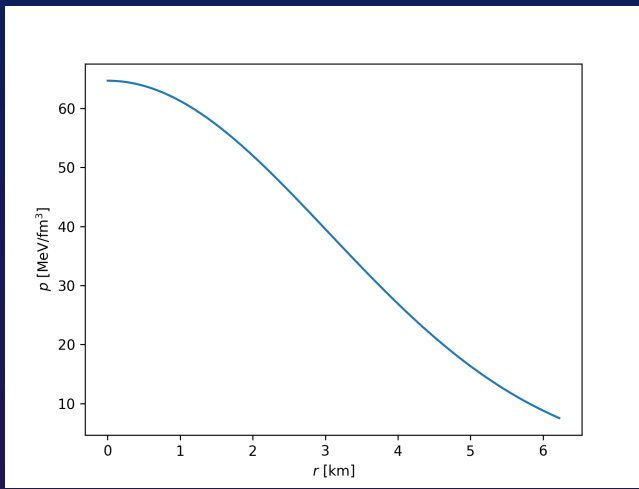
$$M(r) = \int_0^r dr' 4\pi r'^2 \varepsilon(r')$$

- **Equation of state (EoS)** $\varepsilon(p)$
- **Boundary conditions:**
 - **Pressure at the surface** $p(R) = 0$
(in practice $p(R) = p_{\min} = 10^{-5} \text{ km}^{-2}$)
 - **Central energy density** ε_C



M-R diagrams





Kaluza-Klein theory

- Originally: unified field theory of gravity and EM
- **One extra curled up (compactified) dimension**
- **Geometrization of forces, e.g.:**
 - **Newtonian gravity** – gravity is a **force**
 - **General relativity** – gravity arises from the **curved geometry of spacetime**
- We describe the **mass spectrum of nucleons** with the help of an extra dimension



Oskar Klein (1894-1977)



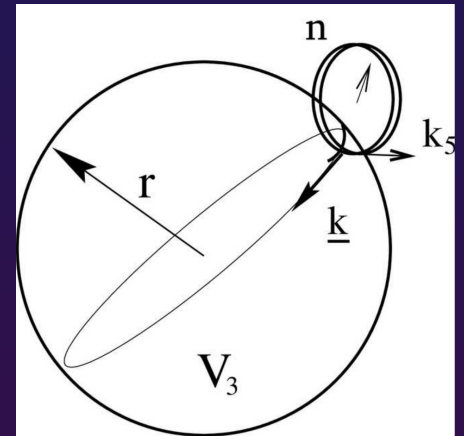
Theodor Kaluza (1885-1954)

Treatment in $1+3+1_c$ dimensions

- Assume one extra **compactified spatial dimension** with size R_C
- At **each point** in ordinary 3D space **particles** with enough energy **can move into it**
 - **3D**: particles with **different masses**
 - **3+1_cD**: **one particle** but with **different quantized momenta** in the extra dimension

$$E_5 = \sqrt{\underline{k}^2 + \left(\frac{n}{R_C}\right)^2} + m^2 = \sqrt{\underline{k}^2 + \bar{m}^2}$$

$$\bar{m}^2 = \left(\frac{n}{R_C}\right)^2 + m^2$$



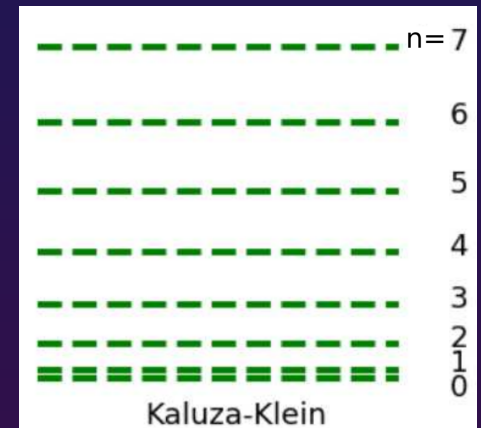
- With the **right** choice of R_C the **mass spectrum** of particles could be **reproduced**

Treatment in 1+3+1_c dimensions

- Assume one extra **compactified spatial dimension** with size R_C
- At **each point** in ordinary 3D space **particles** with enough energy **can move into it**
 - **3D**: particles with **different masses**
 - **3+1_cD**: **one particle** but with **different quantized momenta** in the extra dimension

$$E_5 = \sqrt{\underline{k}^2 + \left(\frac{n}{R_C}\right)^2} + m^2 = \sqrt{\underline{k}^2 + \bar{m}^2}$$

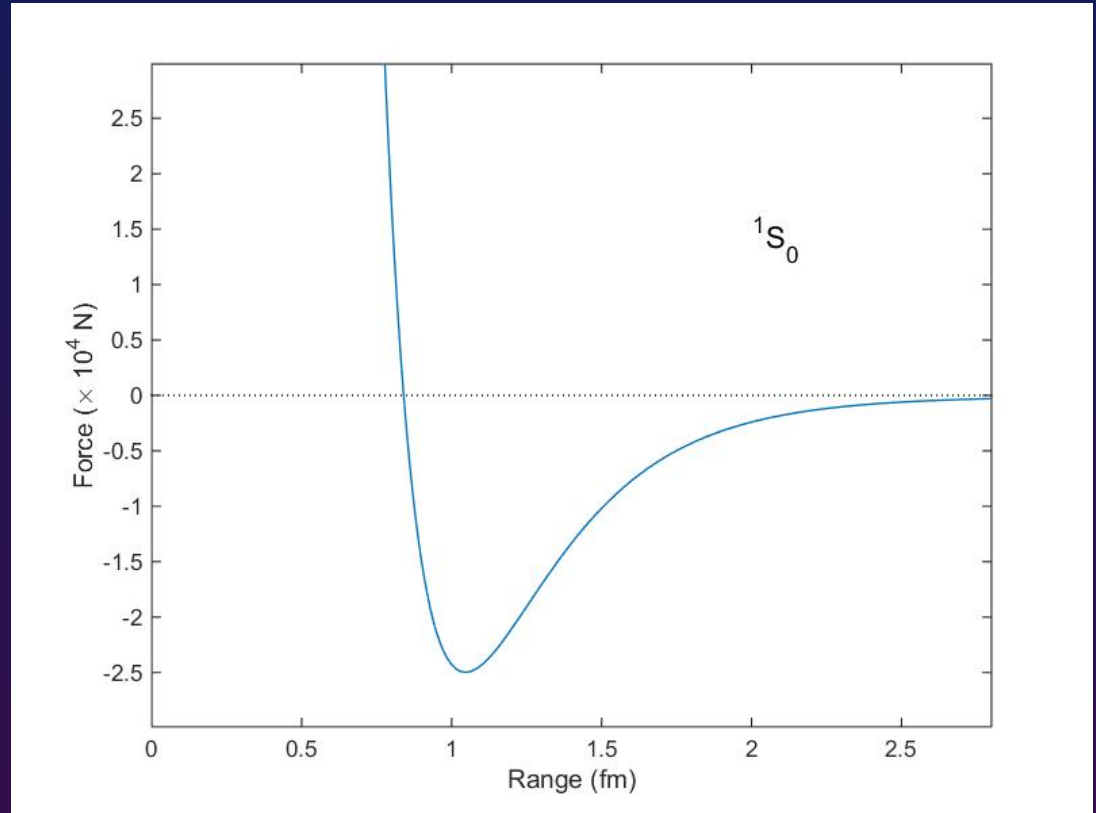
$$\bar{m}^2 = \left(\frac{n}{R_C}\right)^2 + m^2$$



- With the **right** choice of R_C the **mass spectrum** of particles could be **reproduced**

We also need interaction...

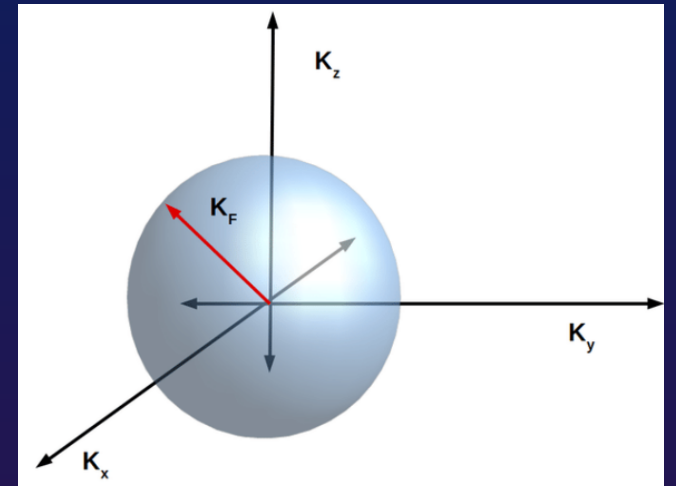
- Pauli exclusion principle is not the only thing that keeps particles apart
- We need to model the **strong force**:
 - **repulsive** at this scale
- Without interaction the maximum mass of stars would become too small $\sim 0.7M_{\odot}$



EoS in 1+4D

- Interacting degenerate Fermi gas
- Potential is a linear function of density:
$$U(n) = \xi n \quad \xi = \text{const}$$
- Thermodynamic potential on $T=0$ MeV

$$\tilde{\Omega} = \frac{-2k_B T V_{(d)}}{h^d} \int \ln \left(1 + e^{\frac{\mu - E(\mathbf{p})}{k_B T}} \right) d^d \mathbf{p}$$



- Extra dimension \longrightarrow calculate with excited mass
- Interaction \longrightarrow chemical potential shifted by $-U(n)$

$$\epsilon(\mu) = \epsilon_0(\mu - U(n)) + \epsilon_{int}$$

$$p(\mu) = p_0(\mu - U(n)) + p_{int}$$

$$n(\mu) = n_0(\mu - U(n))$$

$$\epsilon_{int} = p_{int} = \int U(n) dn = \int \xi n dn = \frac{1}{2} \xi n^2$$

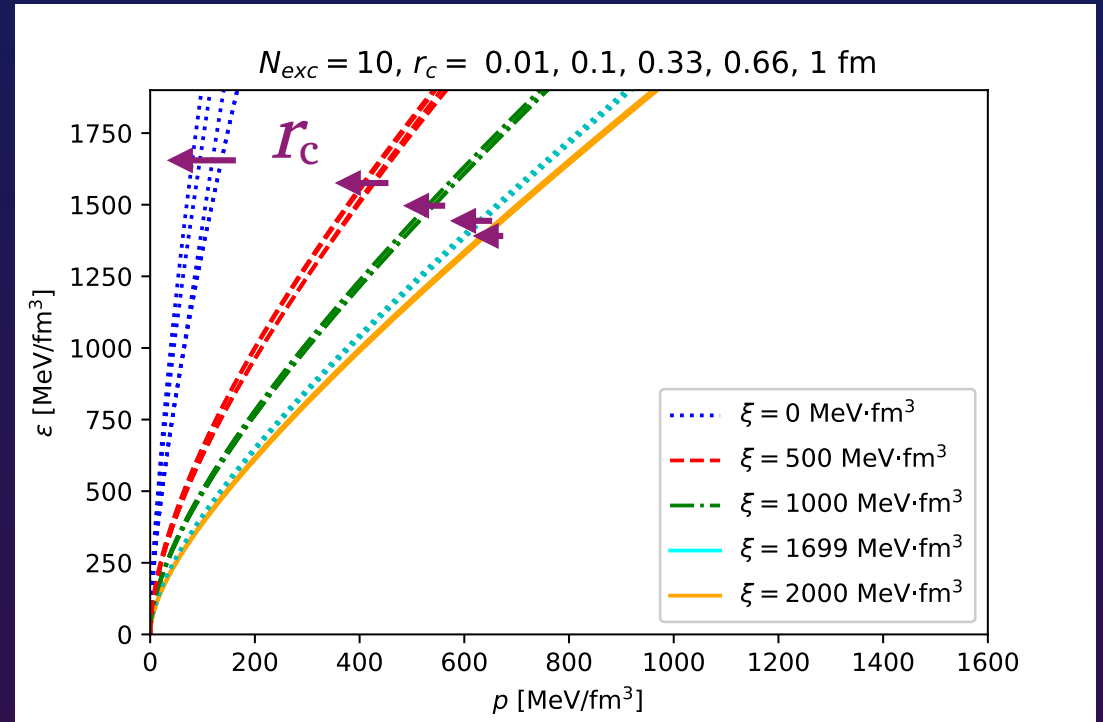
Relativity in 1+4D

- Assume (for **TOV**):
 - ◊ Spherical symmetry
 - ◊ Time-independence
 - ◊ Isotropic relativistic ideal fluid
- Assume (for **extra dimension**):
 - Microscopic
 - ◊ 4D metric does not depend on g_{55}
 - ◊ Causality postulates hold
 - ◊ Full Killing symmetry

$$g_{\mu\nu} = \begin{bmatrix} g_{00} & \cancel{g_{01}} & 0 & 0 & \cancel{g_{05}} \\ \cancel{g_{01}} & g_{11} & 0 & 0 & \cancel{g_{15}} \\ 0 & 0 & g_{22} & 0 & 0 \\ 0 & 0 & 0 & g_{22} \sin^2 \vartheta & 0 \\ \cancel{g_{05}} & \cancel{g_{15}} & 0 & 0 & \boxed{g_{55}} \end{bmatrix}$$

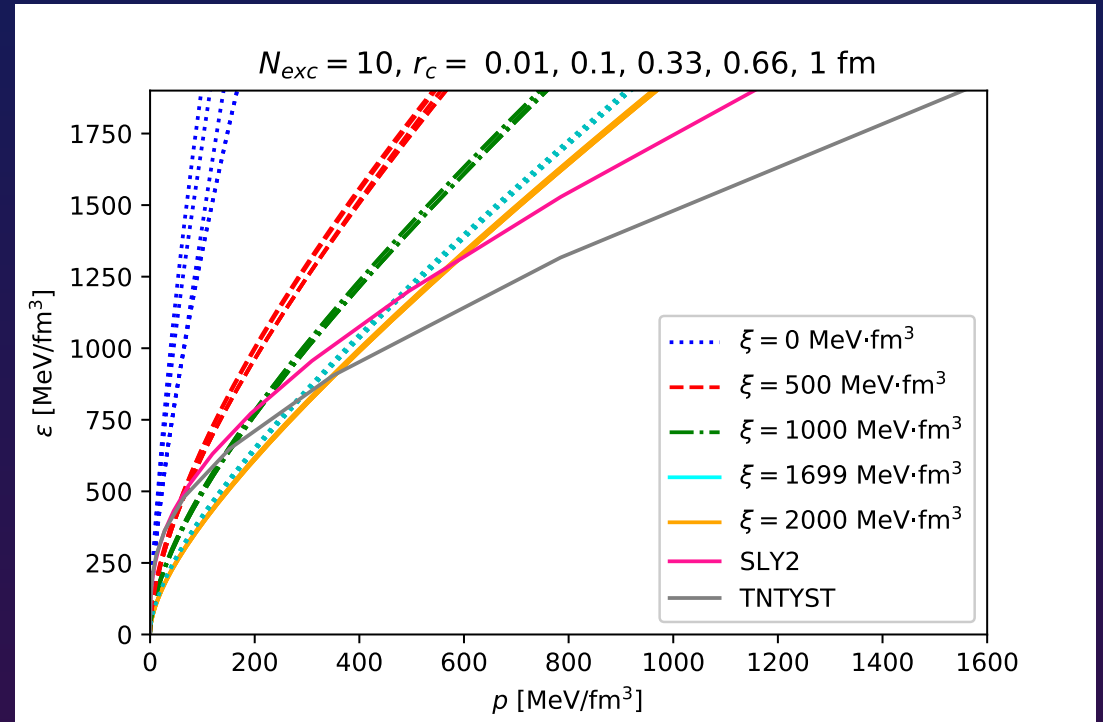
Equation of state

- ξ dependence is much more dominant than r_c
- The bigger ξ , the less important r_c becomes



Equation of state

- ξ dependence is much more dominant than r_c
- The bigger ξ , the less important r_c becomes
- For lower energies small ξ approximates more realistic EoSs
- For high energies a large ξ is a better approximation

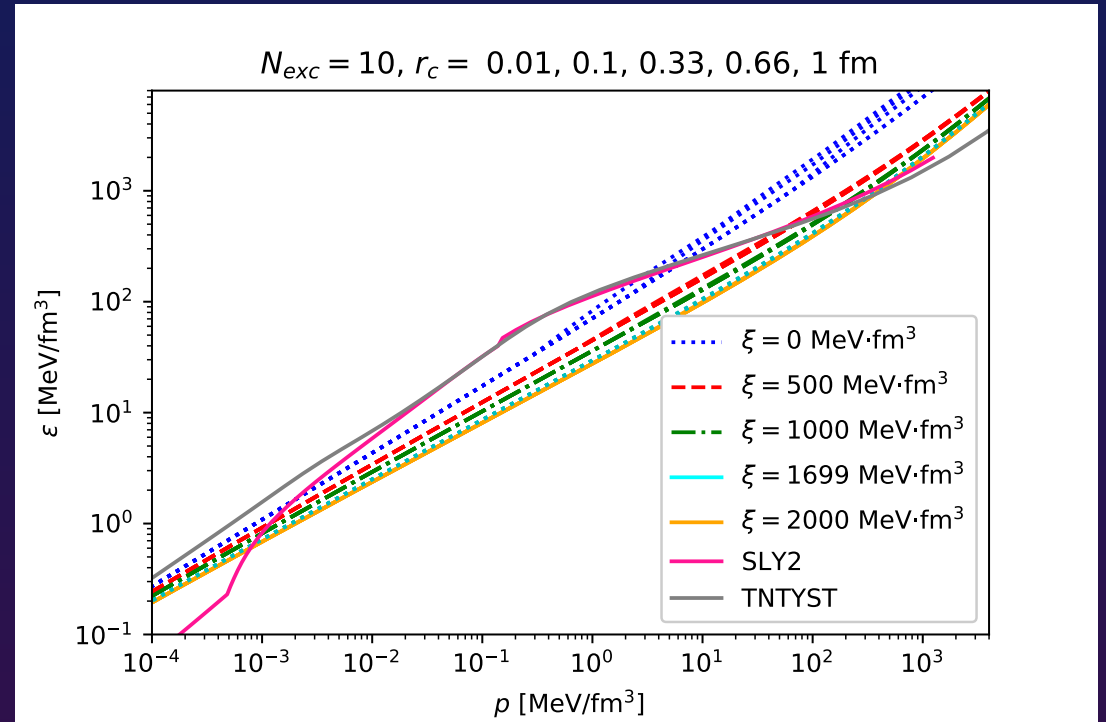


<https://compose.obspm.fr/>

1. H. Togashi, K. Nakazato, Y. Takehara, S. Yamamuro, H. Suzuki and M. Takano, Nucl. Phys. A 961 (2017) 78
1. F. Gulminelli and Ad. R. Raduta, arXiv:1504.04493.
2. E. Chabanat, Ph.D. thesis, University Claude Bernard Lyon-1, Lyon, France, 1995.
3. P. Danielewicz et J. Lee, Nucl. Phys. A818, 36 (2009).

Equation of state

- ξ dependence is much more dominant than r_c
- The bigger ξ , the less important r_c becomes
- For lower energies small ξ approximates more realistic EoSs
- For high energies a large ξ is a better approximation

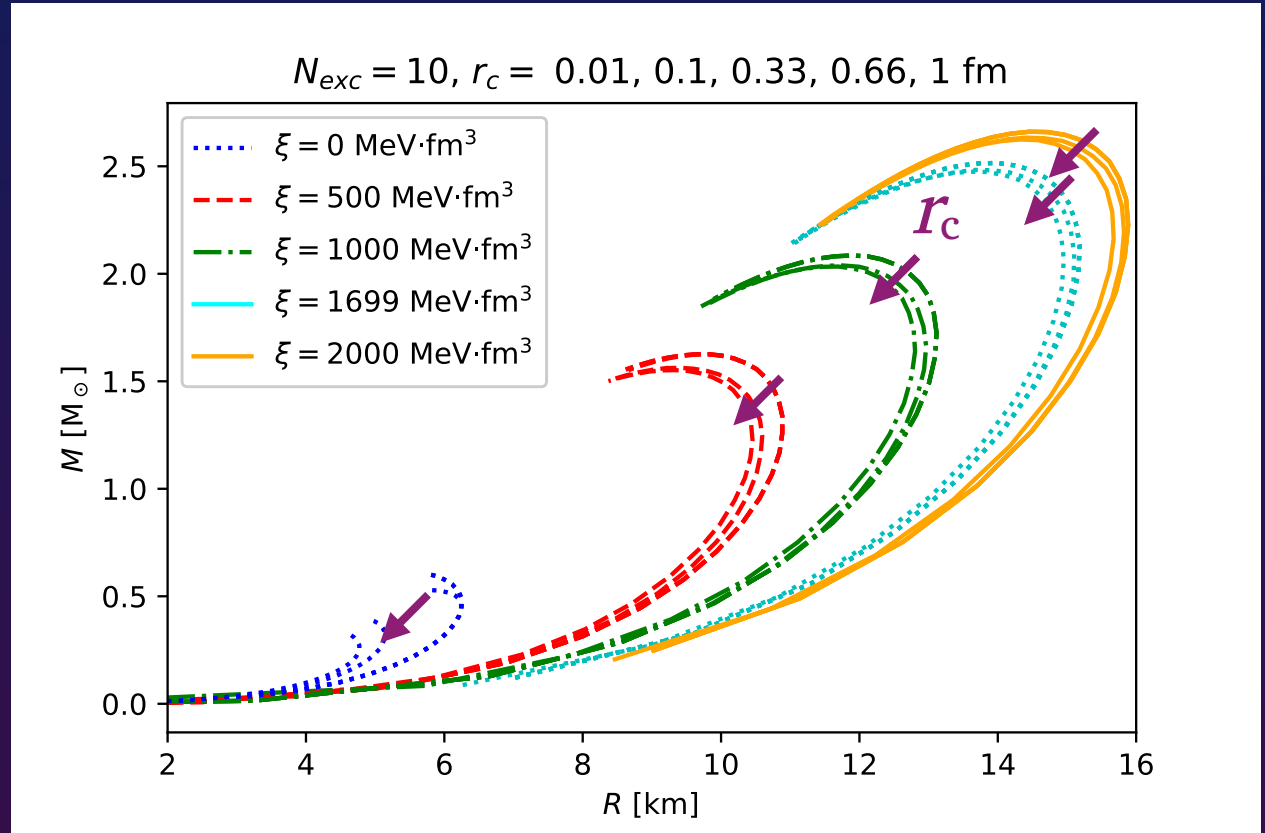


<https://compose.obspm.fr/>

1. H. Togashi, K. Nakazato, Y. Takehara, S. Yamamuro, H. Suzuki and M. Takano, Nucl. Phys. A 961 (2017) 78
1. F. Gulminelli and Ad. R. Raduta, arXiv:1504.04493.
2. E. Chabanat, Ph.D. thesis, University Claude Bernard Lyon-1, Lyon, France, 1995.
3. P. Danielewicz et J. Lee, Nucl. Phys. A818, 36 (2009).

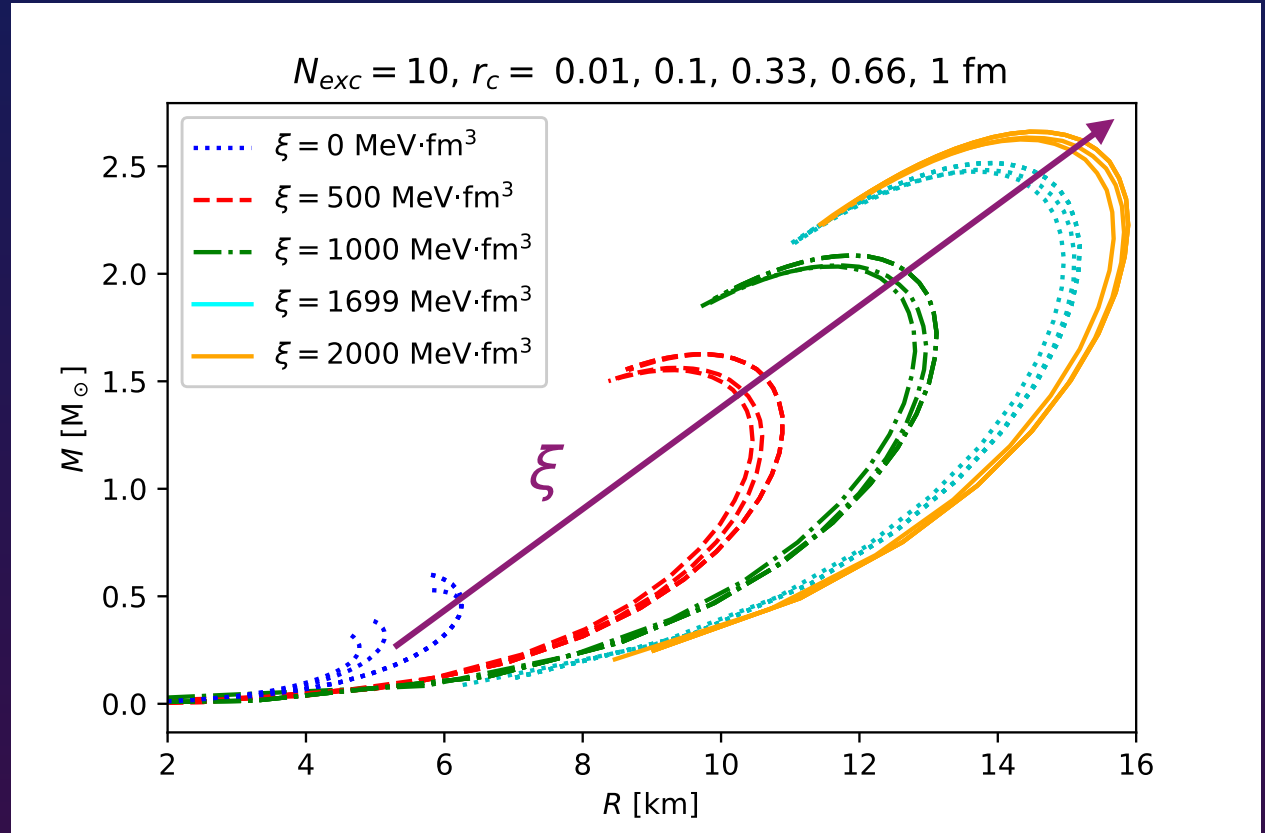
M - R diagrams of the EoS

- ξ dependence is much more dominant than r_c (latter only $\sim 5\%$)
- The bigger ξ , the less important r_c becomes



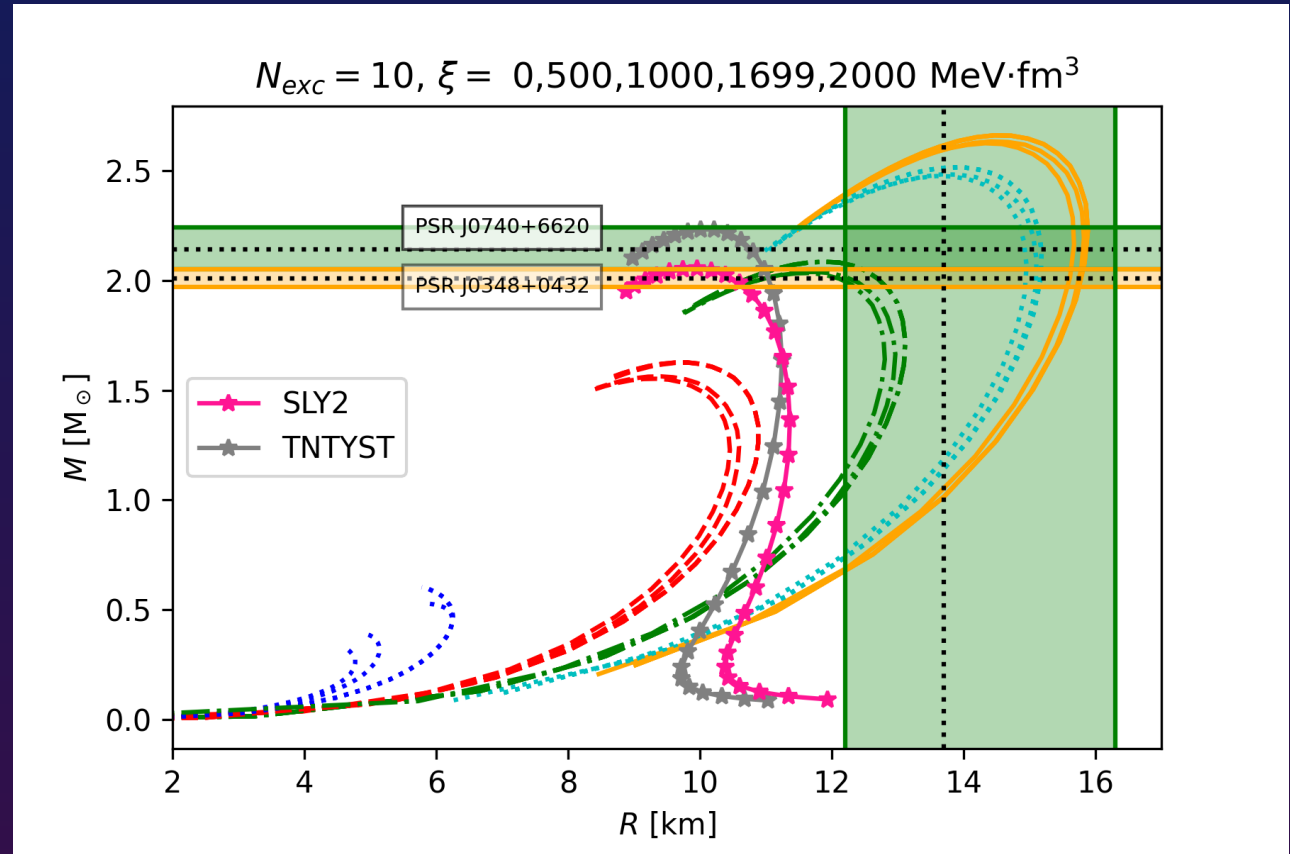
M - R diagrams of the EoS

- ξ dependence is much more dominant than r_c (latter only $\sim 5\%$)
- The bigger ξ , the less important r_c becomes



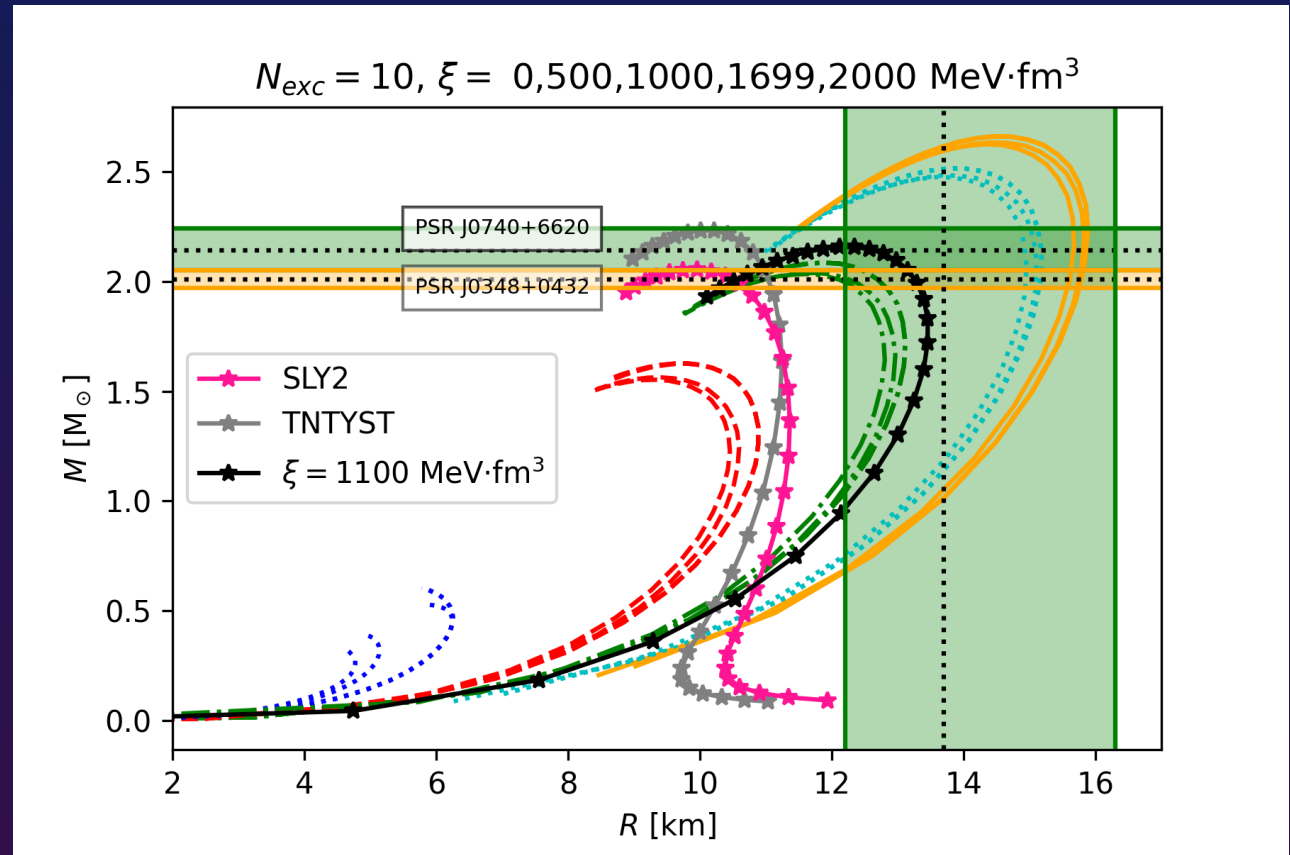
M - R diagrams

- + measurement data
- + 2 more realistic EoSs



M - R diagrams

- + measurement data
- + 2 more realistic EoSs
- + approximation for ξ

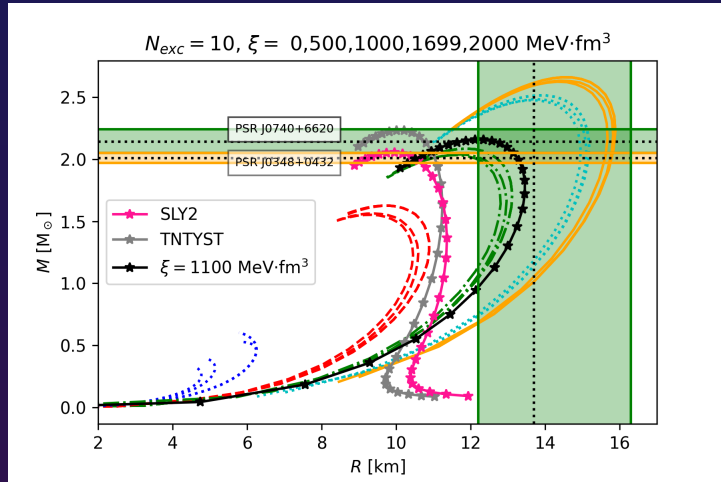


Summary

- Model with the possibility of probing **beyond standard model physics**
- One **extra spatial compactified dimension**



Ordinary **mass** can be described as **quantized 5thD momenta**



- Effective nuclear field theory with **linear repulsive potential**
- It is **possible** to **build** compact stars with **realistic** properties
- **Constraints** on the **size of possible extra dimensions** could be given