

THE TSALLIS-THERMOMETER AS A QGP INDICATOR FOR LARGE AND SMALL COLLISIONAL SYSTEMS

ELTE PARTICLE PHYSICS SEMINAR

GÁBOR BÍRÓ

2021 February 16

WIGNER RESEARCH CENTRE FOR
EÖTVÖS LORÁND **PHYSICS**
UNIVERSITY

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G. Bíró, G.G. Barnaföldi, T.S. Bíró, J. Phys. G, 47.10 (2020), 105002.

G. Bíró, G.G. Barnaföldi, K. Ürmösy, T.S. Bíró, Á. Takács, Entropy, 19(3), (2017), 88

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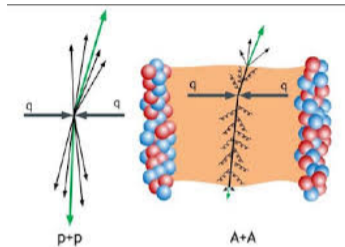
MOTIVATION

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Recap from **RÓBERT VÉRTESI: SCALING PROPERTIES OF JETS IN HIGH-ENERGY PP COLLISIONS:**

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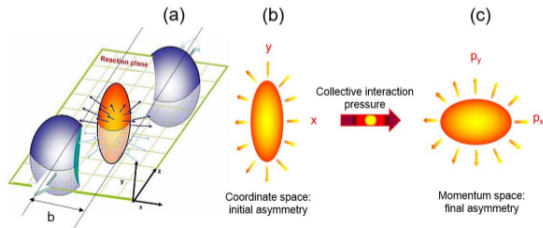
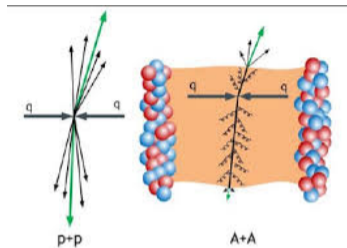
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Collectivity in small systems:

- Observed at high-multiplicity events at LHC (**PRL 112, (2014), 082301**)
- QGP is **not necessary** for collectivity(?)

Jets in high-multiplicity events



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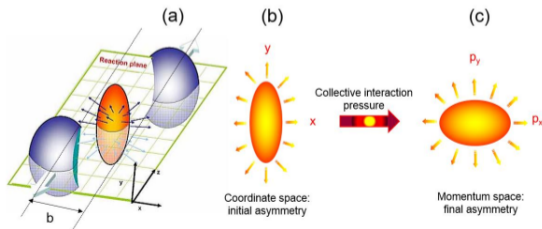
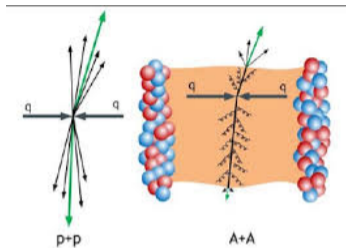
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**"BASIC QUESTION:
CAN WE TURN THE
QGP OFF?"**



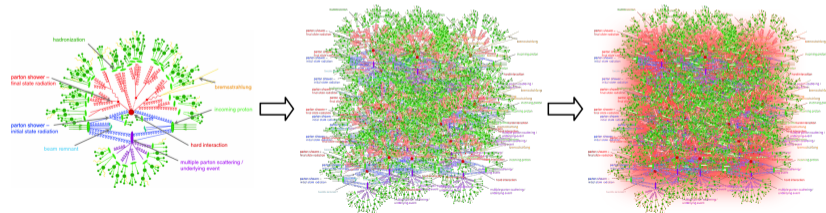
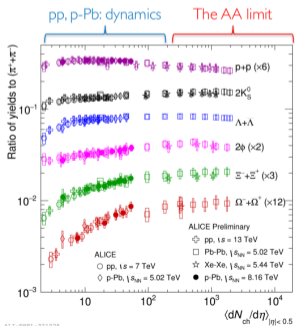
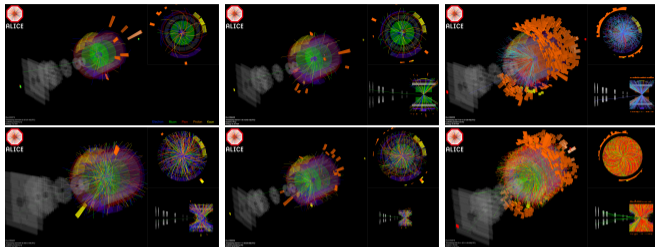
MOTIVATION

Experimental observable:

Ratio of identified hadrons in small to large systems...

...but what is **small**?

Small systems can have **large** multiplicities too...



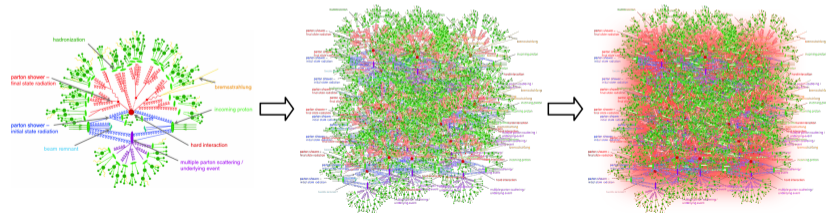
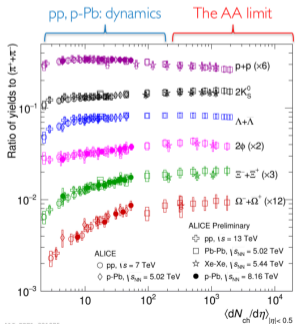
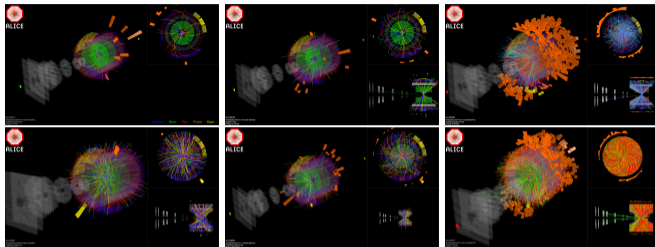
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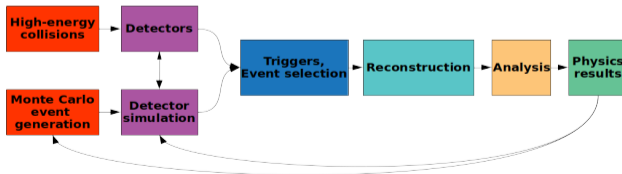
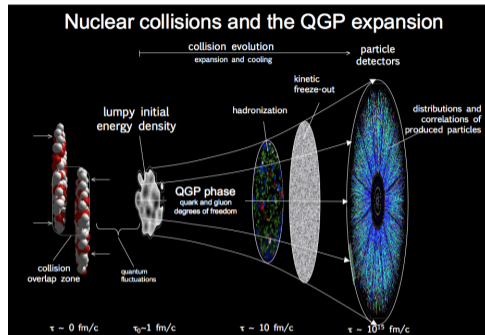
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Where does the quark-gluon plasma start in multiplicity?

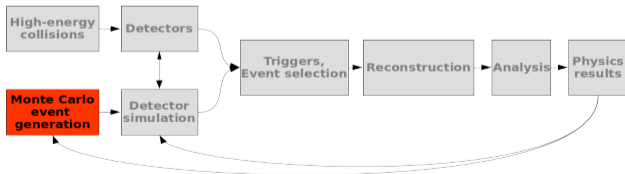
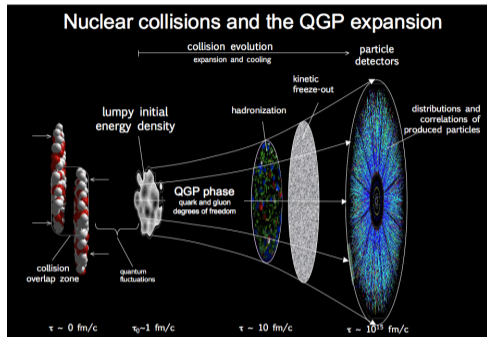


How to connect the **theory** (lattice QCD, ...) with the **experiment** (p_T spectra, multiplicity selection, ...)?



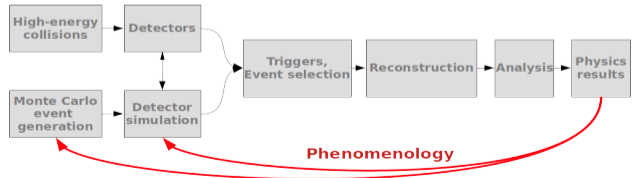
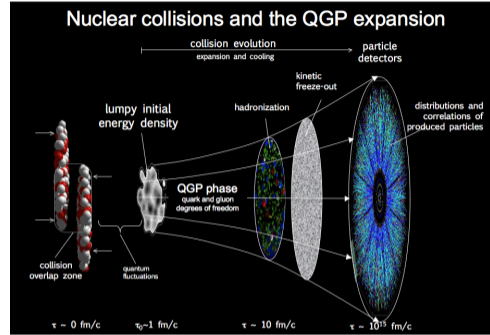
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1. Develop theory to explain the observation (**HIJING++**)



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1. Develop theory to explain the observation (**HIJING++**)
2. Investigate the experimental results with **phenomenological** methods



PRELUDE

The Tsallis – Pareto-type fit functions describe the hadron spectra well – any distribution form could work, but the physical considerations in the background may differ.

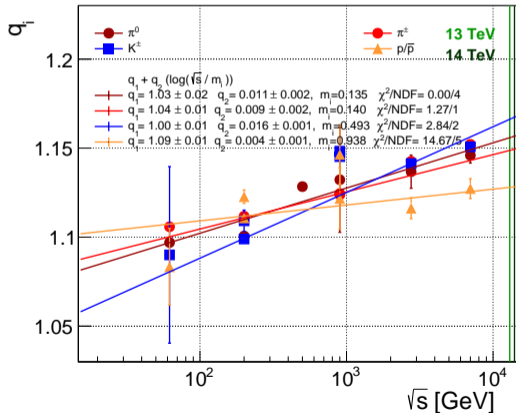
$$\frac{d^2 N}{2\pi dy p_T dp_T} = \begin{cases} A \left(1 + \frac{E}{T}(q-1)\right)^{-\frac{1}{q-1}}, \\ A m_T \left(1 + \frac{E}{T}(q-1)\right)^{-\frac{1}{q-1}}, \\ A \frac{(n-1)(n-2)}{2\pi n T [nT + m(n-2)]} \left(1 + \frac{E}{T}(q-1)\right)^{-\frac{1}{q-1}}, \\ \dots \end{cases}$$

where $n = 1/(q-1)$, and

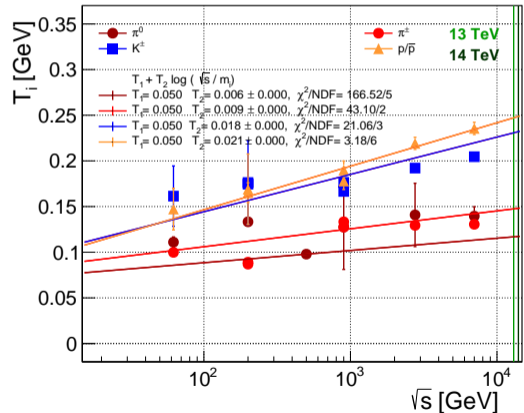
$$E = \begin{cases} p_T, \\ m_T, \\ \gamma (m_T - v p_T), \\ \dots \end{cases}$$

The fitted parameters depend on the center-of-mass energy:

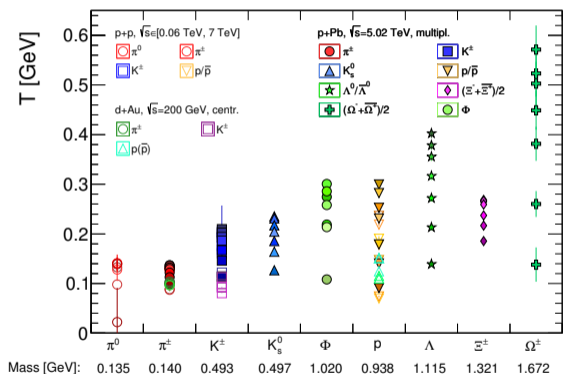
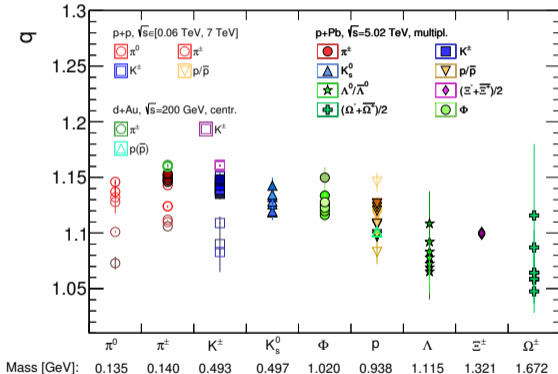
$$q(\sqrt{s}) = q_1 + q_2 \log(\sqrt{s}/m)$$



$$T(\sqrt{s}) = T_1 + T_2 \log(\sqrt{s}/m)$$

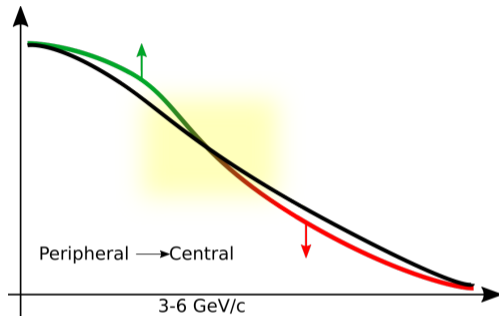


The fitted parameters present a strong mass hierarchy:



GOALS

Non-extensive statistics – summary:



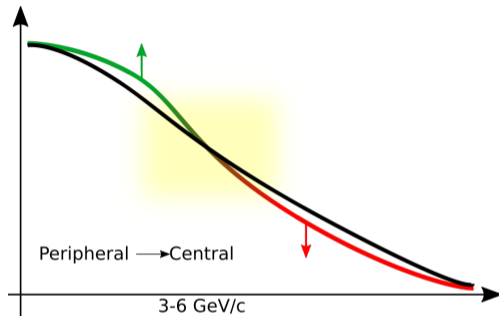
$$\frac{d^2N}{2\pi p_T dp_T dy} = A m_T \left[1 + \frac{q-1}{T} (m_T - m) \right]^{-\frac{q}{q-1}}$$

Non-extensive statistics – summary:

 q -entropy:

$$S_q = \frac{1}{q-1} \left(1 - \sum_{i=1}^W p_i^q \right)$$

$$\lim_{q \rightarrow 1} S_q = S_{BG}$$



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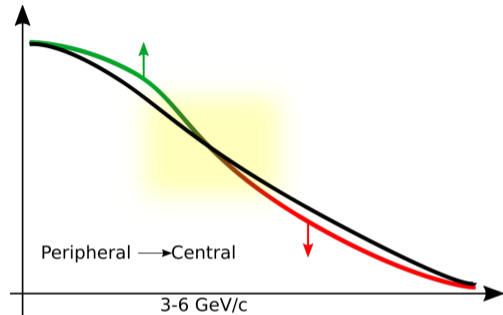
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Thermodynamical consistency:

$$P = Ts + \mu n - \varepsilon$$

$$P = g \int \frac{d^3 p}{(2\pi)^3} T f \quad s = g \int \frac{d^3 p}{(2\pi)^3} \left[\frac{E - \mu}{T} f^q + f \right]$$

$$N = nV = gV \int \frac{d^3 p}{(2\pi)^3} f^q \quad \varepsilon = g \int \frac{d^3 p}{(2\pi)^3} E f^q$$



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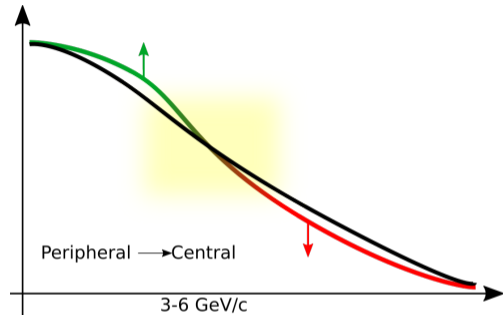
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Final size effects:

$$T = \frac{E}{\langle n \rangle} \quad q = 1 - \frac{1}{\langle n \rangle} + \frac{\Delta n^2}{\langle n \rangle^2}$$

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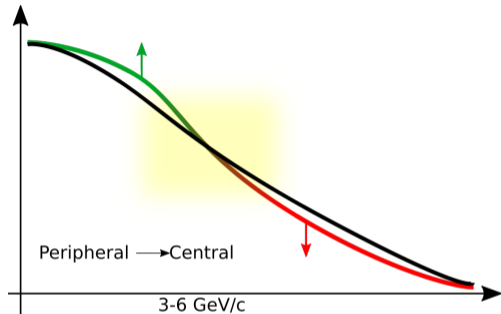
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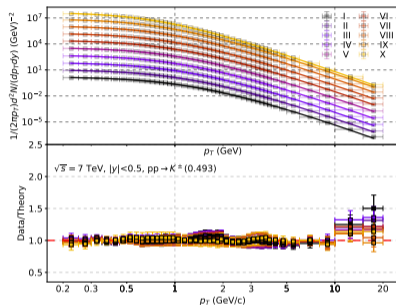
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Strong indication for multiplicity (system size) dependency:

$$\frac{\langle dN_{ch}/d\eta \rangle}{\langle N_{part} \rangle / 2} \propto \begin{cases} s_{NN}^{0.15} & \text{for AA,} \\ s^{0.11} & \text{for pp.} \end{cases}$$

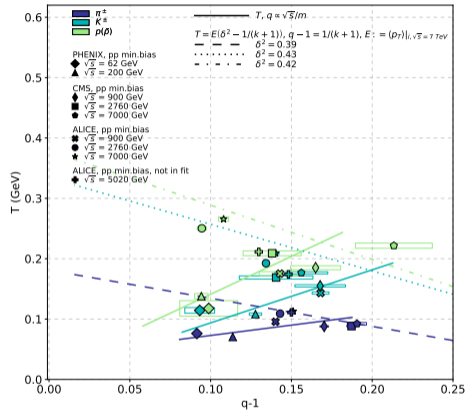
What is the relation with the earlier observations?

"Min. bias" pp: \sim low multiplicity

1. solid lines: \sqrt{s} dependency from earlier
2. $\sqrt{s} \sim$ multiplicity \sim NBD

$$q \sim NBD \quad (q - 1 = 1/(k + 1))$$

3. $T = E(\delta^2 - (q - 1)), E := \langle p_T \rangle |_{\sqrt{s}=7 TeV}$



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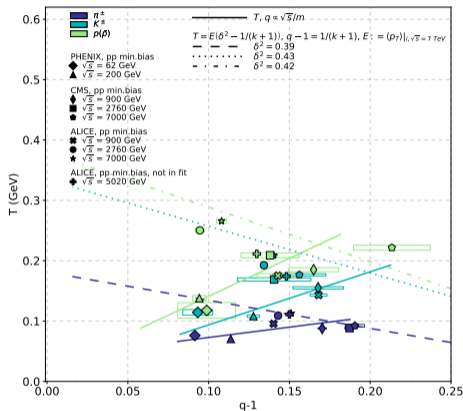
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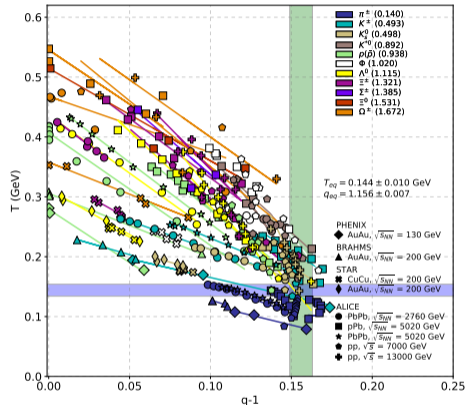


System size suggests an opposite trend!

The approach:

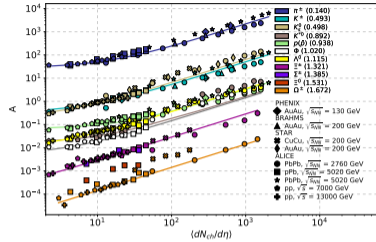
Map the thermodynamically consistent non-extensive parameter space of the available experimental data and compare it with theoretical QCD calculations

- 11 identified hadron species: from π^\pm to Ω
- Various collision systems: proton-proton, proton-nucleus, nucleus-nucleus
- Wide range of multiplicities: $2.2 \leq \langle dN_{ch}/d\eta \rangle \leq 2047$
- Wide range of CM energies: $130 \leq \sqrt{s_{NN}} \leq 13000$ GeV
- **More than 30** published experimental datasets



**Goal: calibrate the
Tsallis-thermometer**

RESULTS

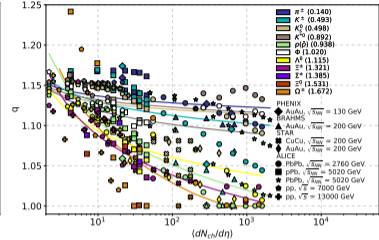
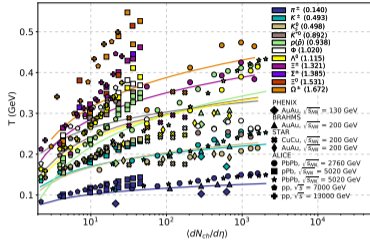


Parametrizations:

$$A = A_0 + A_1 \ln \frac{\sqrt{s_{NN}}}{m} + A_2 \langle dN_{ch}/d\eta \rangle$$

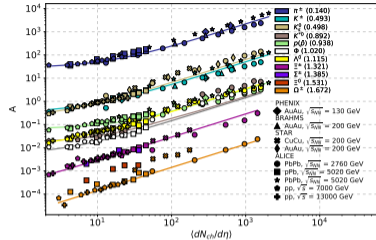
$$T = T_0 + T_1 \ln \frac{\sqrt{s_{NN}}}{m} + T_2 \ln \ln \langle dN_{ch}/d\eta \rangle$$

$$q = q_0 + q_1 \ln \frac{\sqrt{s_{NN}}}{m} + q_2 \ln \ln \langle dN_{ch}/d\eta \rangle$$



1. The **A**, **q** and **T** parameters characterize the collision
2. Strong **grouping**: $T_{eq} \approx 0.144$ GeV, $q_{eq} \approx 1.156$

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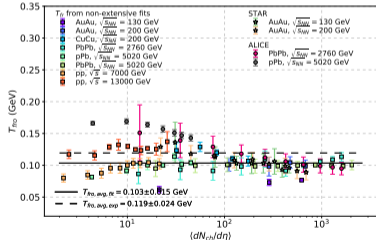
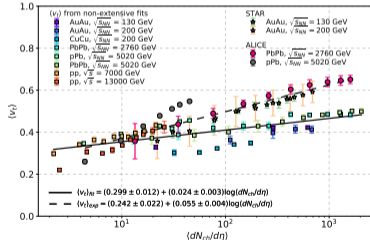
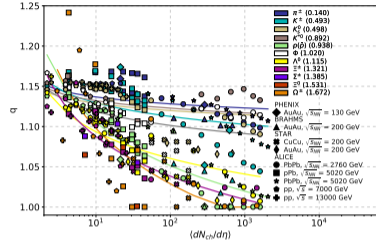
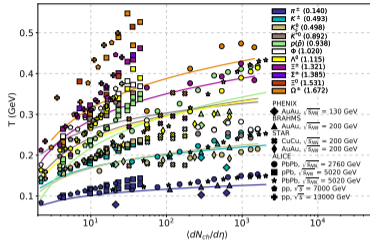
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Radial flow:

$$T = T_{fro} + m \langle u_t \rangle^2$$

$$\langle v_t \rangle = \frac{\langle u_t \rangle}{\sqrt{1 + \langle u_t \rangle^2}}$$



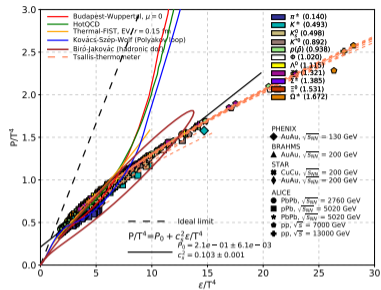
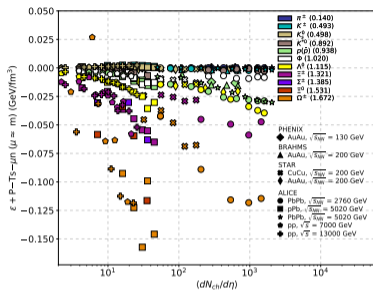
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2. Strong **grouping**: $T_{eq} \approx 0.144$ GeV, $q_{eq} \approx 1.156$

3. **Test**: results are comparable with experiments (**Phys. Rev. C 83 (2011), 064903**)

Thermodynamical consistency: ✓

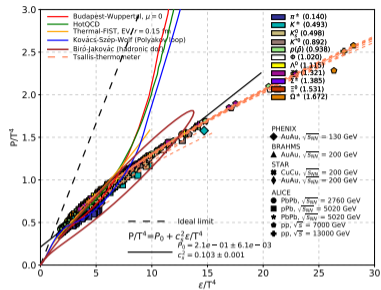
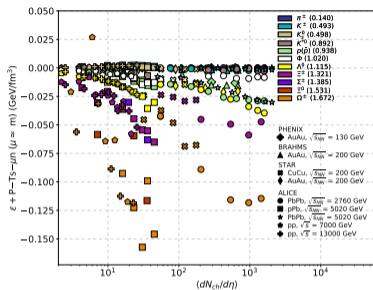
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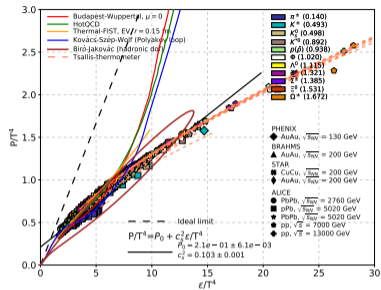
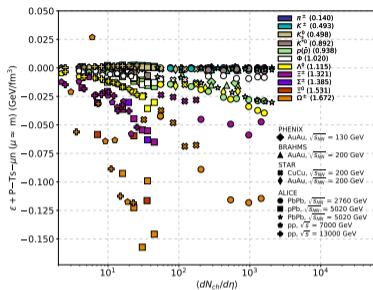
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Comparison of the thermodynamical variables with theoretical calculations (with **one giant leap for theory**):

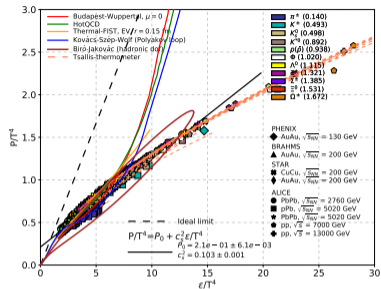
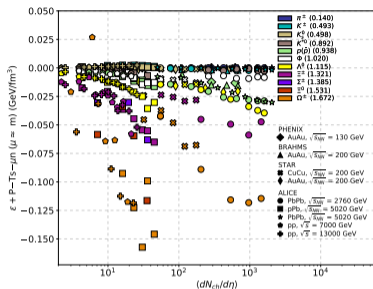


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Comparison of the thermodynamical variables with theoretical calculations (with **one giant leap for theory**):

Hadronic dof \Leftrightarrow **partonic dof**



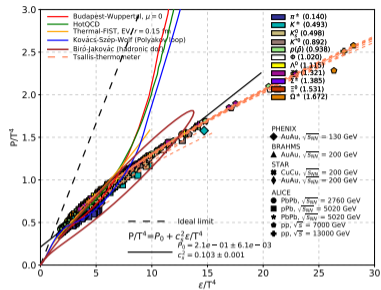
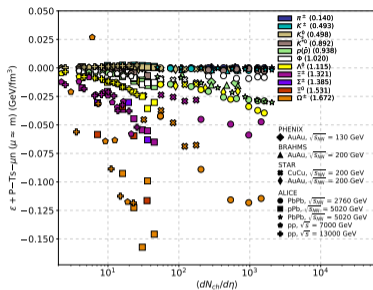
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$$P = Ts + \mu n - \varepsilon$$

Comparison of the thermodynamical variables with theoretical calculations (with **one giant leap for theory**):

Hadronic dof \Leftrightarrow partonic dof

Interpretation of the grouping phenomenon in the $T - (q - 1)$ parameter space:



Thermodynamical consistency: ✓

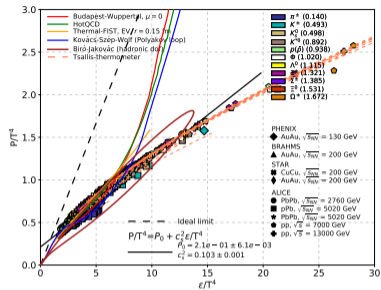
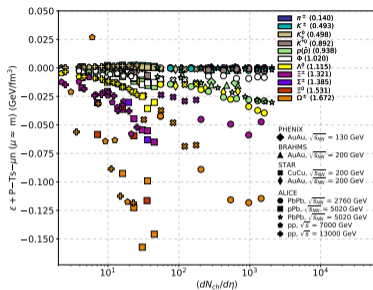
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Comparison of the thermodynamical variables with theoretical calculations (with **one giant leap for theory**):

Hadronic dof \Leftrightarrow partonic dof

Interpretation of the grouping phenomenon in the $T - (q - 1)$ parameter space:

1. Overlapping region with theoretical calculations \rightarrow **presence of hot QCD matter** just before the hadronization?



Thermodynamical consistency: ✓

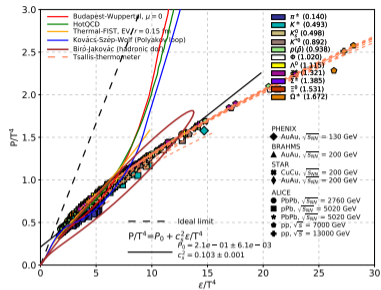
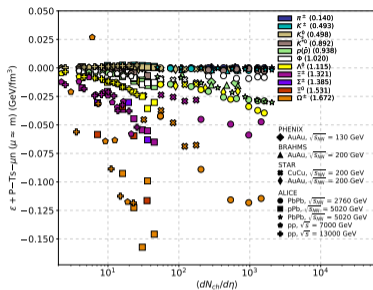
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1. Overlapping region with theoretical calculations \rightarrow **presence of hot QCD matter** just before the hadronization?
2. Hadron spectra of colliding systems with $T \approx 0.144$ GeV and $q \approx 1.156$: originates from a previous **quark-gluon plasma** state



Thermodynamical consistency: ✓

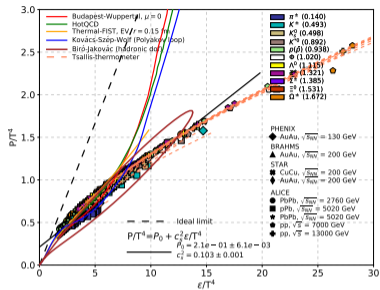
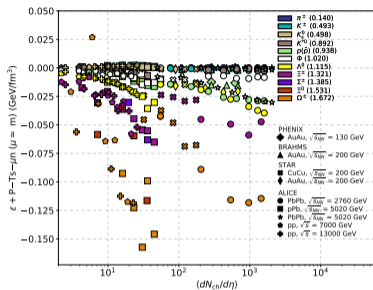
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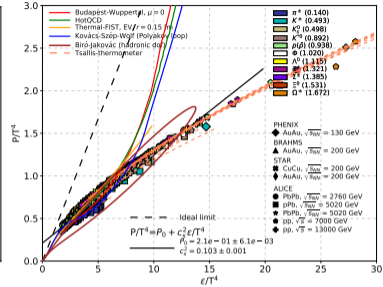
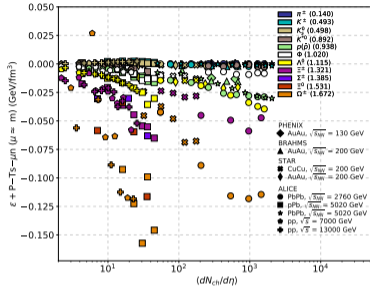
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With the **parametrizations**: \sqrt{s} and $\langle dN_{ch}/d\eta \rangle$ regions:



SUMMARY

- Consistent non-extensive analysis of a **very large set** of experimental data
- New results are in agreement with earlier studies
- $q \neq 1$ for all hadron spectra: dependency on the size of the collisional system through **multiplicity** fluctuations
- **Various checks** of the non-extensive framework
- Grouping of the **T** and **q** parameters
- **Comparison** with theoretical QCD calculations
- **Tsallis-thermometer**: final state hadrons may originate from a previously present strongly interacting QCD matter at event multiplicities as low as $\langle dN_{ch}/d\eta \rangle \sim 100$

SUPPORT

The research is supported by: OTKA K120660, K123815, K135515, THOR COST CA15213, Hungarian-Chinese 12 CN-1-2012-0016, MOST 2014DFG02050, 2019-2.1.11-TÉT-2019-00050 Tét, Wigner HAS-OBOR-CCNU, ÚNKP-17-3.

Thank you for your attention!

BACKUP

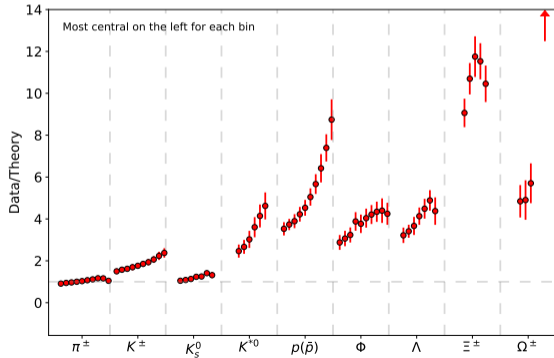
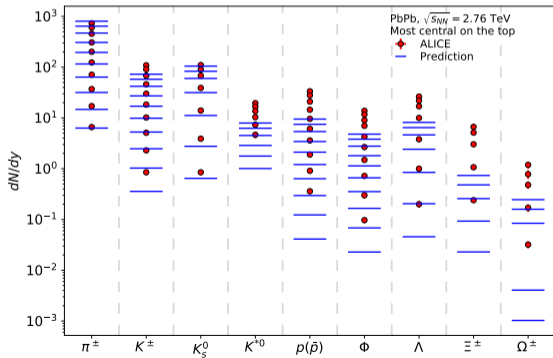
EXPERIMENTAL DATA

System, $\sqrt{s_{NN}}$ (GeV)	η or y	Hadron	Mult. classes	p_T range (GeV/ c)	
AuAu, 130	$ \eta < 0,35$	π^\pm	3, [21,3; 622]	[0,25; 2,2]	
		K^\pm		[0,45; 1,65]	
CuCu, 200	$ y < 0,5$	$p(\bar{p})$	5, [32; 175]	[0,55; 3,42]	
		K_s^0		[0,5; 9,0]	
		Λ^0		[0,5; 7,0]	
		Ξ^\pm		[0,7; 6,0]	
		Ω^\pm		[1,0; 4,5]	
		Φ		[0,45; 4,5]	
AuAu, 200	$ y < 0,2$	π^\pm	3, [111; 680]	[0,2; 2,0]	
		K^\pm		[0,4; 2,0]	
PbPb, 2760	$ y < 0,5$	$p(\bar{p})$	5, [27; 680]	[0,3; 3,0]	
		K_s^0		[0,5; 9,0]	
			Λ^0		[0,5; 8,0]
	$ y < 0,5$	10, [13,4; 1601]	π^\pm	[0,1; 3,0]	
			K^\pm	[0,2; 3,0]	
			K_s^0	[0,4; 12,0]	
			K^{*0}	[0,3; 20,0]	
			$p(\bar{p})$	[0,3; 4,6]	
			Λ^0	[0,6; 12,0]	
			Φ	[0,5; 21,0]	
Ξ^\pm			[0,6; 8,0]		
pPb, 5020	$-0,5 < y < 0,0$	π^\pm	7, [4,3; 45]	[0,1; 20,0]	
		K^\pm		[0,2; 20,0]	
		K^{*0}		[0,0; 16,0]	
		$p(\bar{p})$		[0,35; 20,0]	
		Φ		[0,4; 20,0]	
		Ξ^0		[0,8; 8,0]	

System, $\sqrt{s_{NN}}$ (GeV)	η or y	Hadron	Mult. classes	p_T range (GeV/ c)
pPb, 5020	$-0,5 < y < 0,0$	π^\pm	7, [4,3; 45]	[0,1; 20,0]
		Σ^\pm		[1,0; 6,0]
		Ξ^\pm		[0,6; 7,2]
	$0,0 < y < 0,5$	Ω^\pm	7, [4,3; 45]	[0,8; 5,0]
		π^\pm		[0,1; 3,0]
		K^\pm		[0,2; 2,4]
PbPb, 5020	$ y < 0,5$	K_s^0	10, [19,5; 2047]	[0,0; 8,0]
		$p(\bar{p})$		[0,3; 4,0]
		Λ^0		[0,6; 8,0]
		π^\pm		[0,1; 10,0]
		K^\pm		[0,1; 10,0]
		$p(\bar{p})$		[0,1; 10,0]
pp, 7000	$ y < 0,5$	π^\pm	10, [2,2; 21,3]	[0,1; 20,0]
		K^\pm		[0,2; 20,0]
		K_s^0		[0,0; 12,0]
		K^{*0}		[0,0; 10,0]
		$p(\bar{p})$		[0,3; 20,0]
		Φ		[0,4; 10,0]
		Λ^0		[0,4; 8,0]
		Ξ^\pm		[0,6; 6,5]
		Ω^\pm		[0,9; 5,5]
pp, 13000	$ y < 0,5$	K_s^0	10, [2,52; 25,72]	[0,0; 12,0]
		Λ^0		[0,4; 8,0]
		Ξ^\pm		[0,6; 6,5]
		Ω^\pm		[0,9; 5,5]

HADRON YIELDS FROM THE PARAMETRIZATIONS

$$\left. \frac{dN}{dy} \right|_{y=0} = 2\pi A T \left[\frac{(2-q)m^2 + 2mT + 2T^2}{(2-q)(3-2q)} \right] \left[1 + \frac{q-1}{T} m \right]^{-\frac{1}{q-1}}$$



THERMODYNAMICAL QUANTITIES

